

**FIGURE 21-38**

Solar radiation reaching the earth's atmosphere and the total solar irradiance.

on the path of the infrared radiation coming from the earth, and thus slows down the cooling process at night. In areas with clear skies such as deserts, there is a large swing between the daytime and nighttime temperatures because of the absence of such barriers for infrared radiation.

## 21-6 ■ ATMOSPHERIC AND SOLAR RADIATION

The sun is our primary source of energy. The energy coming off the sun, called *solar energy*, reaches us in the form of electromagnetic waves after experiencing considerable interactions with the atmosphere. The radiation energy emitted or reflected by the constituents of the atmosphere form the *atmospheric radiation*. Here we give an overview of the solar and atmospheric radiation because of their importance and relevance to daily life. Also, our familiarity with solar energy makes it an effective tool in developing a better understanding for some of the new concepts introduced earlier. Detailed treatment of this exciting subject can be found in numerous books devoted to this topic.

The *sun* is a nearly spherical body that has a diameter of  $D \approx 1.39 \times 10^9$  m and a mass of  $m \approx 2 \times 10^{30}$  kg and is located at a mean distance of  $L = 1.50 \times 10^{11}$  m from the earth. It emits radiation energy continuously at a rate of  $E_{\text{sun}} \approx 3.8 \times 10^{26}$  W. Less than a billionth of this energy (about  $1.7 \times 10^{17}$  W) strikes the earth, which is sufficient to keep the earth warm and to maintain life through the photosynthesis process. The energy of the sun is due to the continuous *fusion* reaction during which two hydrogen atoms fuse to form one atom of helium. Therefore, the sun is essentially a *nuclear reactor*, with temperatures as high as 40,000,000 K in its core region. The temperature drops to about 5800 K in the outer region of the sun, called the convective zone, as a result of the dissipation of this energy by radiation.

The solar energy reaching the earth's atmosphere is called the **total solar irradiance**  $G_s$ , whose value is

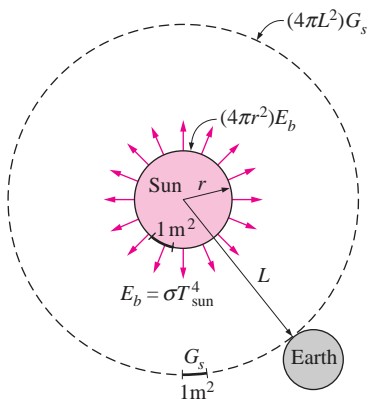
$$G_s = 1373 \text{ W/m}^2 \quad (21-49)$$

The total solar irradiance (also called the **solar constant**) represents *the rate at which solar energy is incident on a surface normal to the sun's rays at the outer edge of the atmosphere when the earth is at its mean distance from the sun* (Fig. 21-38).

The value of the total solar irradiance can be used to estimate the effective surface temperature of the sun from the requirement that

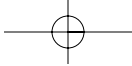
$$(4\pi L^2)G_s = (4\pi r^2)\sigma T_{\text{sun}}^4 \quad (21-50)$$

where  $L$  is the mean distance between the sun's center and the earth and  $r$  is the radius of the sun. The left-hand side of this equation represents the total solar energy passing through a spherical surface whose radius is the mean earth-sun distance, and the right-hand side represents the total energy that leaves the sun's outer surface. The conservation of energy principle requires that these two quantities be equal to each other, since the solar energy experiences no attenuation (or enhancement) on its way through the vacuum (Fig. 21-39). The **effective surface temperature** of the sun is determined from Eq. 21-50 to be  $T_{\text{sun}} = 5780$  K. That is, the sun can be treated as a



**FIGURE 21-39**

The total solar energy passing through concentric spheres remains constant, but the energy falling per unit area decreases with increasing radius.



blackbody at a temperature of 5780 K. This is also confirmed by the measurements of the spectral distribution of the solar radiation just outside the atmosphere plotted in Fig. 21–40, which shows only small deviations from the idealized blackbody behavior.

The spectral distribution of solar radiation on the ground plotted in Fig. 21–40 shows that the solar radiation undergoes considerable *attenuation* as it passes through the atmosphere as a result of *absorption* and *scattering*. About 99 percent of the atmosphere is contained within a distance of 30 km from the earth's surface. The several dips on the spectral distribution of radiation on the earth's surface are due to *absorption* by the gases  $O_2$ ,  $O_3$  (ozone),  $H_2O$ , and  $CO_2$ . Absorption by *oxygen* occurs in a narrow band about  $\lambda = 0.76 \mu\text{m}$ . The *ozone* absorbs *ultraviolet* radiation at wavelengths below  $0.3 \mu\text{m}$  almost completely, and radiation in the range  $0.3\text{--}0.4 \mu\text{m}$  considerably. Thus, the ozone layer in the upper regions of the atmosphere protects biological systems on earth from harmful ultraviolet radiation. In turn, we must protect the ozone layer from the destructive chemicals commonly used as refrigerants, cleaning agents, and propellants in aerosol cans. The use of these chemicals is now banned in many countries. The ozone gas also absorbs some radiation in the visible range. Absorption in the infrared region is dominated by *water vapor* and *carbon dioxide*. The dust particles and other pollutants in the atmosphere also absorb radiation at various wavelengths.

As a result of these absorptions, the solar energy reaching the *earth's surface* is weakened considerably, to about  $950 \text{ W/m}^2$  on a clear day and much less on cloudy or smoggy days. Also, practically all of the solar radiation reaching the earth's surface falls in the wavelength band from  $0.3$  to  $2.5 \mu\text{m}$ .

Another mechanism that attenuates solar radiation as it passes through the atmosphere is *scattering* or *reflection* by air molecules and the many other kinds of particles such as dust, smog, and water droplets suspended in the atmosphere. Scattering is mainly governed by the size of the particle relative to the wavelength of radiation. The oxygen and nitrogen molecules primarily scatter radiation at very short wavelengths, comparable to the size of the molecules themselves. Therefore, radiation at wavelengths corresponding to violet and blue colors is scattered the most. This molecular scattering in all directions is what gives the sky its bluish color. The same phenomenon is responsible for red sunrises and sunsets. Early in the morning and late in the afternoon, the sun's rays pass through a greater thickness of the atmosphere than they do at midday, when the sun is at the top. Therefore, the violet and blue colors of the light encounter a greater number of molecules by the time they reach the earth's surface, and thus a greater fraction of them are scattered (Fig. 21–41). Consequently, the light that reaches the earth's surface consists primarily of colors corresponding to longer wavelengths such as red, orange, and yellow. The clouds appear in reddish-orange color during sunrise and sunset because the light they reflect is reddish-orange at those times. For the same reason, a red traffic light is visible from a longer distance than is a green light under the same circumstances.

The solar energy incident on a surface on earth is considered to consist of *direct* and *diffuse* parts. The part of solar radiation that reaches the earth's surface without being scattered or absorbed by the atmosphere is called **direct solar radiation**  $G_D$ . The scattered radiation is assumed to reach the earth's surface uniformly from all directions and is called **diffuse solar radiation**  $G_d$ .

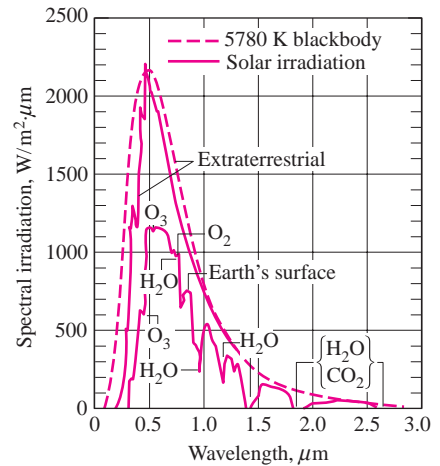


FIGURE 21–40

Spectral distribution of solar radiation just outside the atmosphere, at the surface of the earth on a typical day, and comparison with blackbody radiation at 5780 K.

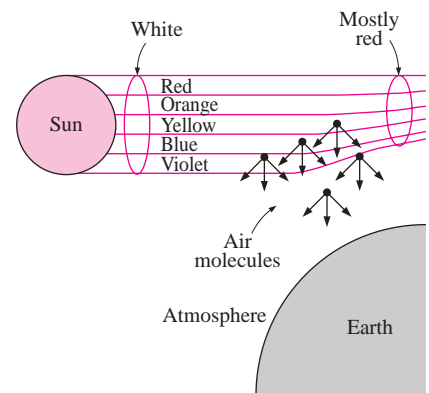
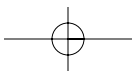
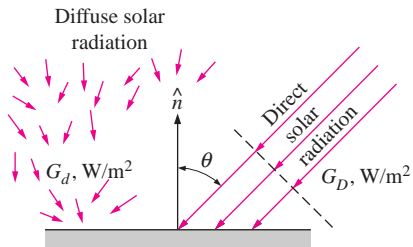
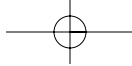


FIGURE 21–41

Air molecules scatter blue light much more than they do red light. At sunset, the light travels through a thicker layer of atmosphere, which removes much of the blue from the natural light, allowing the red to dominate.





**FIGURE 21-42**  
The direct and diffuse radiation incident on a horizontal surface at the earth's surface.

Then the *total solar energy* incident on the unit area of a *horizontal surface* on the ground is (Fig. 21-42)

$$G_{\text{solar}} = G_D \cos \theta + G_d \quad (\text{W/m}^2) \quad (21-51)$$

where  $\theta$  is the angle of incidence of direct solar radiation (the angle that the sun's rays make with the normal of the surface). The diffuse radiation varies from about 10 percent of the total radiation on a clear day to nearly 100 percent on a totally cloudy day.

The gas molecules and the suspended particles in the atmosphere *emit radiation* as well as absorbing it. The atmospheric emission is primarily due to the  $\text{CO}_2$  and  $\text{H}_2\text{O}$  molecules and is concentrated in the regions from 5 to 8  $\mu\text{m}$  and above 13  $\mu\text{m}$ . Although this emission is far from resembling the distribution of radiation from a blackbody, it is found convenient in radiation calculations to treat the atmosphere as a blackbody at some lower fictitious temperature that emits an equivalent amount of radiation energy. This fictitious temperature is called the **effective sky temperature**  $T_{\text{sky}}$ . Then the radiation emission from the atmosphere to the earth's surface is expressed as

$$G_{\text{sky}} = \sigma T_{\text{sky}}^4 \quad (\text{W/m}^2) \quad (21-52)$$

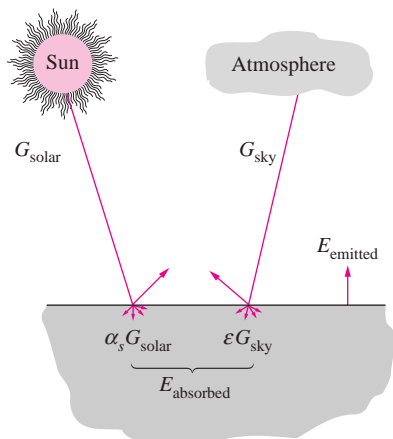
The value of  $T_{\text{sky}}$  depends on the atmospheric conditions. It ranges from about 230 K for cold, clear-sky conditions to about 285 K for warm, cloudy-sky conditions.

Note that the effective sky temperature does not deviate much from the room temperature. Thus, in the light of Kirchhoff's law, we can take the absorptivity of a surface to be equal to its emissivity at room temperature,  $\alpha = \varepsilon$ . Then the sky radiation absorbed by a surface can be expressed as

$$E_{\text{sky, absorbed}} = \alpha G_{\text{sky}} = \alpha \sigma T_{\text{sky}}^4 = \varepsilon \sigma T_{\text{sky}}^4 \quad (\text{W/m}^2) \quad (21-53)$$

The net rate of radiation heat transfer to a surface exposed to solar and atmospheric radiation is determined from an energy balance (Fig. 21-43):

$$\begin{aligned} \dot{q}_{\text{net, rad}} &= \sum E_{\text{absorbed}} - \sum E_{\text{emitted}} \\ &= E_{\text{solar, absorbed}} + E_{\text{sky, absorbed}} - E_{\text{emitted}} \\ &= \alpha_s G_{\text{solar}} + \varepsilon \sigma T_{\text{sky}}^4 - \varepsilon \sigma T_s^4 \\ &= \alpha_s G_{\text{solar}} + \varepsilon \sigma (T_{\text{sky}}^4 - T_s^4) \quad (\text{W/m}^2) \end{aligned} \quad (21-54)$$

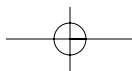


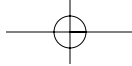
**FIGURE 21-43**  
Radiation interactions of a surface exposed to solar and atmospheric radiation.

where  $T_s$  is the temperature of the surface in K and  $\varepsilon$  is its emissivity at room temperature. A positive result for  $\dot{q}_{\text{net, rad}}$  indicates a radiation heat gain by the surface and a negative result indicates a heat loss.

The absorption and emission of radiation by the *elementary gases* such as  $\text{H}_2$ ,  $\text{O}_2$ , and  $\text{N}_2$  at moderate temperatures are negligible, and a medium filled with these gases can be treated as a *vacuum* in radiation analysis. The absorption and emission of gases with *larger molecules* such as  $\text{H}_2\text{O}$  and  $\text{CO}_2$ , however, can be *significant* and may need to be considered when considerable amounts of such gases are present in a medium. For example, a 1-m-thick layer of water vapor at 1 atm pressure and 100°C emits more than 50 percent of the energy that a blackbody would emit at the same temperature.

In solar energy applications, the spectral distribution of incident solar radiation is very different than the spectral distribution of emitted radiation by





the surfaces, since the former is concentrated in the short-wavelength region and the latter in the infrared region. Therefore, the radiation properties of surfaces will be quite different for the incident and emitted radiation, and the surfaces cannot be assumed to be gray. Instead, the surfaces are assumed to have two sets of properties: one for solar radiation and another for infrared radiation at room temperature. Table 21–3 lists the *emissivity*  $\varepsilon$  and the *solar absorptivity*  $\alpha_s$  of the surfaces of some common materials. Surfaces that are intended to *collect solar energy*, such as the absorber surfaces of solar collectors, are desired to have high  $\alpha_s$  but low  $\varepsilon$  values to maximize the absorption of solar radiation and to minimize the emission of radiation. Surfaces that are intended to *remain cool* under the sun, such as the outer surfaces of fuel tanks and refrigerator trucks, are desired to have just the opposite properties. Surfaces are often given the desired properties by coating them with thin layers of *selective* materials. A surface can be kept cool, for example, by simply painting it white.

We close this section by pointing out that what we call *renewable energy* is usually nothing more than the manifestation of solar energy in different forms. Such energy sources include wind energy, hydroelectric power, ocean thermal energy, ocean wave energy, and wood. For example, no hydroelectric power plant can generate electricity year after year unless the water evaporates by absorbing solar energy and comes back as a rainfall to replenish the water source (Fig. 21–44). Although solar energy is sufficient to meet the entire energy needs of the world, currently it is not economical to do so because of the low concentration of solar energy on earth and the high capital cost of harnessing it.

### EXAMPLE 21–5 Selective Absorber and Reflective Surfaces

Consider a surface exposed to solar radiation. At a given time, the direct and diffuse components of solar radiation are  $G_D = 400$  and  $G_d = 300$  W/m<sup>2</sup>, and the direct radiation makes a 20° angle with the normal of the surface. The surface temperature is observed to be 320 K at that time. Assuming an effective sky temperature of 260 K, determine the net rate of radiation heat transfer for these cases (Fig. 21–45):

- $\alpha_s = 0.9$  and  $\varepsilon = 0.9$  (gray absorber surface)
- $\alpha_s = 0.1$  and  $\varepsilon = 0.1$  (gray reflector surface)
- $\alpha_s = 0.9$  and  $\varepsilon = 0.1$  (selective absorber surface)
- $\alpha_s = 0.1$  and  $\varepsilon = 0.9$  (selective reflector surface)

**SOLUTION** A surface is exposed to solar and sky radiation. The net rate of radiation heat transfer is to be determined for four different combinations of emissivities and solar absorptivities.

**Analysis** The total solar energy incident on the surface is

$$\begin{aligned} G_{\text{solar}} &= G_D \cos \theta + G_d \\ &= (400 \text{ W/m}^2) \cos 20^\circ + (300 \text{ W/m}^2) \\ &= 676 \text{ W/m}^2 \end{aligned}$$

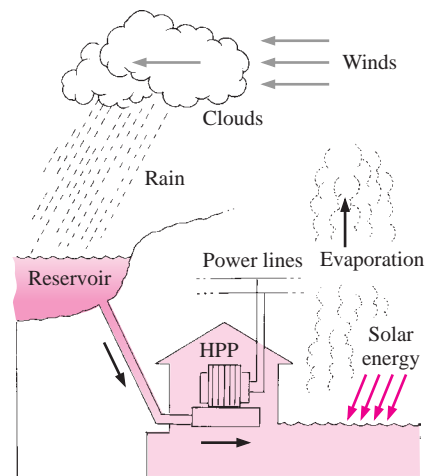
Then the net rate of radiation heat transfer for each of the four cases is determined from:

$$\dot{q}_{\text{net, rad}} = \alpha_s G_{\text{solar}} + \varepsilon \sigma (T_{\text{sky}}^4 - T_s^4)$$

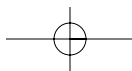
**TABLE 21–3**

Comparison of the solar absorptivity  $\alpha_s$  of some surfaces with their emissivity  $\varepsilon$  at room temperature

Surface	$\alpha_s$	$\varepsilon$
Aluminum		
Polished	0.09	0.03
Anodized	0.14	0.84
Foil	0.15	0.05
Copper		
Polished	0.18	0.03
Tarnished	0.65	0.75
Stainless steel		
Polished	0.37	0.60
Dull	0.50	0.21
Plated metals		
Black nickel oxide	0.92	0.08
Black chrome	0.87	0.09
Concrete	0.60	0.88
White marble	0.46	0.95
Red brick	0.63	0.93
Asphalt	0.90	0.90
Black paint	0.97	0.97
White paint	0.14	0.93
Snow	0.28	0.97
Human skin		
(Caucasian)	0.62	0.97


**FIGURE 21–44**

The cycle that water undergoes in a hydroelectric power plant.



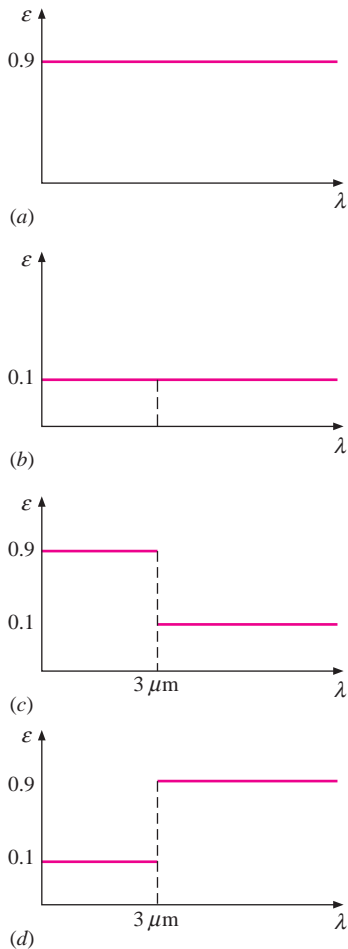


FIGURE 21-45

Graphical representation of the spectral emissivities of the four surfaces considered in Example 21-5.

## SUMMARY

Radiation propagates in the form of electromagnetic waves. The *frequency*  $\nu$  and *wavelength*  $\lambda$  of electromagnetic waves in a medium are related by  $\lambda = c/\nu$ , where  $c$  is the speed of propagation in that medium. All matter whose temperature is above absolute zero continuously emits *thermal radiation* as a result of vibrational and rotational motions of molecules, atoms, and electrons of a substance. Temperature is a measure of the strength of these activities at the microscopic level.

A *blackbody* is defined as a *perfect emitter and absorber of radiation*. At a specified temperature and wavelength, no surface can emit more energy than a blackbody. A blackbody absorbs *all* incident radiation, regardless of wavelength and direction. The radiation energy emitted by a blackbody per unit

(a)  $\alpha_s = 0.9$  and  $\varepsilon = 0.9$  (gray absorber surface):

$$\dot{q}_{\text{net, rad}} = 0.9(676 \text{ W/m}^2) + 0.9(5.67 \times 10^{-8} \text{ W/m}^2 \cdot \text{K}^4)[(260 \text{ K})^4 - (320 \text{ K})^4] \\ = \mathbf{307 \text{ W/m}^2}$$

(b)  $\alpha_s = 0.1$  and  $\varepsilon = 0.1$  (gray reflector surface):

$$\dot{q}_{\text{net, rad}} = 0.1(676 \text{ W/m}^2) + 0.1(5.67 \times 10^{-8} \text{ W/m}^2 \cdot \text{K}^4)[(260 \text{ K})^4 - (320 \text{ K})^4] \\ = \mathbf{34 \text{ W/m}^2}$$

(c)  $\alpha_s = 0.9$  and  $\varepsilon = 0.1$  (selective absorber surface):

$$\dot{q}_{\text{net, rad}} = 0.9(676 \text{ W/m}^2) + 0.1(5.67 \times 10^{-8} \text{ W/m}^2 \cdot \text{K}^4)[(260 \text{ K})^4 - (320 \text{ K})^4] \\ = \mathbf{575 \text{ W/m}^2}$$

(d)  $\alpha_s = 0.1$  and  $\varepsilon = 0.9$  (selective reflector surface):

$$\dot{q}_{\text{net, rad}} = 0.1(676 \text{ W/m}^2) + 0.9(5.67 \times 10^{-8} \text{ W/m}^2 \cdot \text{K}^4)[(260 \text{ K})^4 - (320 \text{ K})^4] \\ = \mathbf{-234 \text{ W/m}^2}$$

**Discussion** Note that the surface of an ordinary gray material of high absorptivity gains heat at a rate of  $307 \text{ W/m}^2$ . The amount of heat gain increases to  $575 \text{ W/m}^2$  when the surface is coated with a selective material that has the same absorptivity for solar radiation but a low emissivity for infrared radiation. Also note that the surface of an ordinary gray material of high reflectivity still gains heat at a rate of  $34 \text{ W/m}^2$ . When the surface is coated with a selective material that has the same reflectivity for solar radiation but a high emissivity for infrared radiation, the surface loses  $234 \text{ W/m}^2$  instead. Therefore, the temperature of the surface will decrease when a selective reflector surface is used.

time and per unit surface area is called the *blackbody emissive power*  $E_b$  and is expressed by the *Stefan-Boltzmann law* as

$$E_b(T) = \sigma T^4$$

where  $\sigma = 5.670 \times 10^{-8} \text{ W/m}^2 \cdot \text{K}^4$  is the *Stefan-Boltzmann constant* and  $T$  is the absolute temperature of the surface in K. At any specified temperature, the spectral blackbody emissive power  $E_{b\lambda}$  increases with wavelength, reaches a peak, and then decreases with increasing wavelength. The wavelength at which the peak occurs for a specified temperature is given by *Wien's displacement law* as

$$(\lambda T)_{\text{max power}} = 2897.8 \mu\text{m} \cdot \text{K}$$